

**SIMULATING EXTREME ANISOTROPY
WITHOUT MESH ALIGNMENT**

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**MHD Working Group, Oct. 28 2001
APS 2001, Long Beach CA**

- Nonlinear simulations of high-temperature plasmas must be able to resolve thermal transport anisotropy associated with the magnetic field direction.
 - The saturation of pressure-driven modes is sensitive to changes in magnetic topology due to parallel thermal conduction.
 - The ratio of parallel to perpendicular thermal conductivities leads to one threshold mechanism for neoclassical tearing modes.
- When the magnetic field is aligned with the grid or when the angle between $\hat{\mathbf{b}}$ and the grid is uniform, the anisotropy is accurately represented by standard techniques.
- A nonlinearly evolving magnetic topology requires more sophisticated approaches:
 - There is curvature in the magnetic field.
 - The topology of islands and stochastic regions is three-dimensional.
 - Even 3D automated mesh refinement schemes would be severely challenged by these conditions; 3D refinement is possible but 3D alignment near a separatrix is not.
- The **NIMROD** code has addressed this challenge by using high-order finite-element basis functions, which represent curvature with or without alignment.
 - The increase in spatial convergence rates with basis function order has been verified.
 - Spatial convergence rates are retained with nonuniform meshing.

A quantitative measure of the numerical error can be determined by a simple test problem.

- The simplest thermal conduction problem is a 2D box with Dirichlet boundary conditions and a source. If the source drives the lowest eigenfunction only,

$$S(x, y) = 2\pi^2 \cos(\pi x) \cos(\pi y)$$

in the domain

$$-0.5 \leq x \leq 0.5, -0.5 \leq y \leq 0.5 ,$$

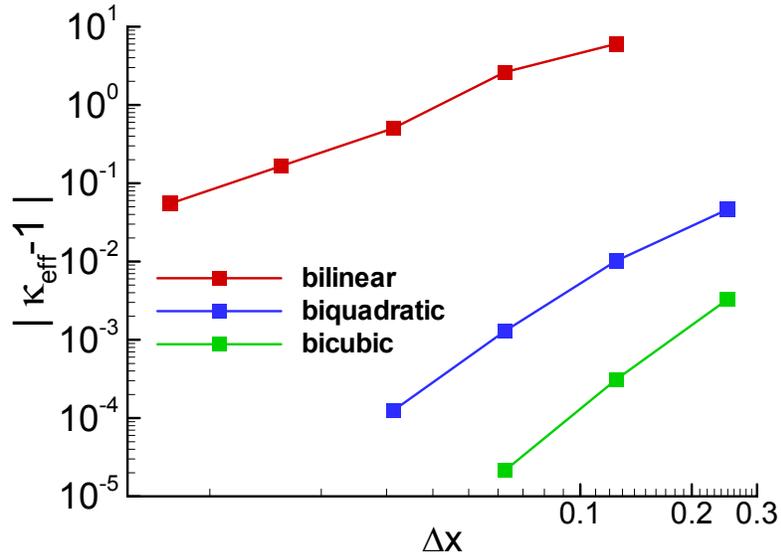
it produces the temperature distribution

$$T(x, y) = \kappa^{-1} \cos(\pi x) \cos(\pi y) .$$

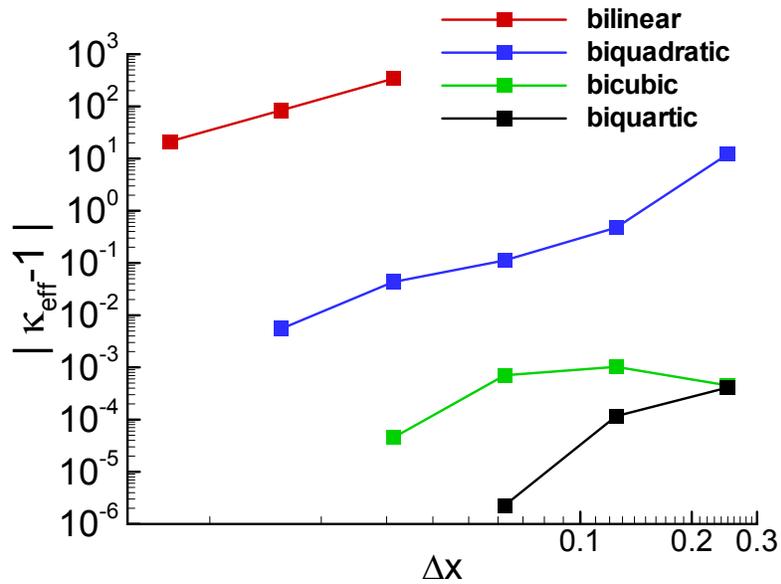
- A numerical test with \mathbf{B} everywhere tangent to this temperature distribution and a uniform rectilinear grid has $\hat{\mathbf{b}}$ severely misaligned with the grid. With $\kappa_{\perp} = 1$ and $\kappa_{\parallel} \gg 1$, the inverse of the resulting $T(0,0)$ is a good measure of the effective κ_{\perp} of the numerical algorithm.
- The magnetic field is created by inducing a perpendicular current density distribution that is proportional to the heat source.

Results from the Anisotropic Diffusion Test

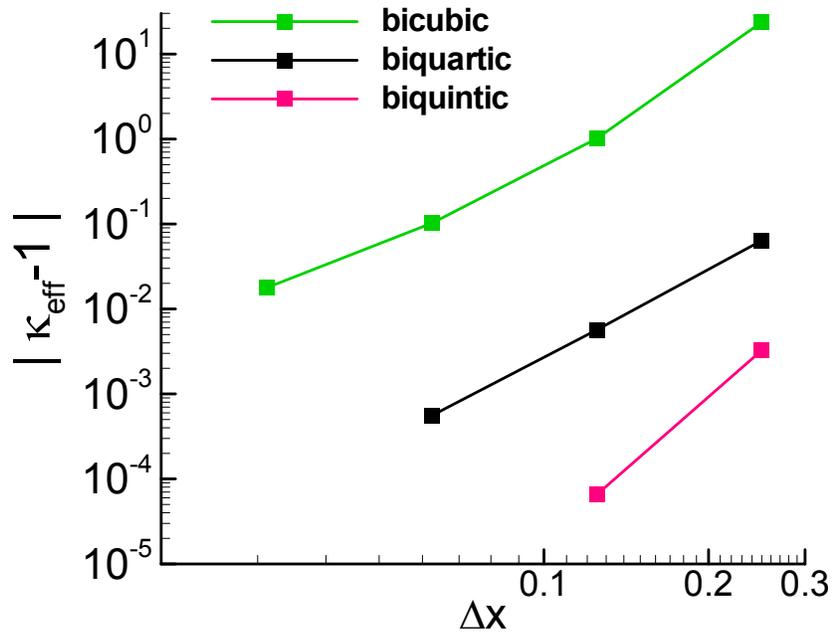
Effective $\kappa_{\text{perp}}^{-1}$ for $\kappa_{\parallel} = 10^3$



Effective $\kappa_{\text{perp}}^{-1}$ for $\kappa_{\parallel} = 10^6$



Effective $\kappa_{\text{perp}}^{-1}$ for $\kappa_{\parallel} = 10^9$



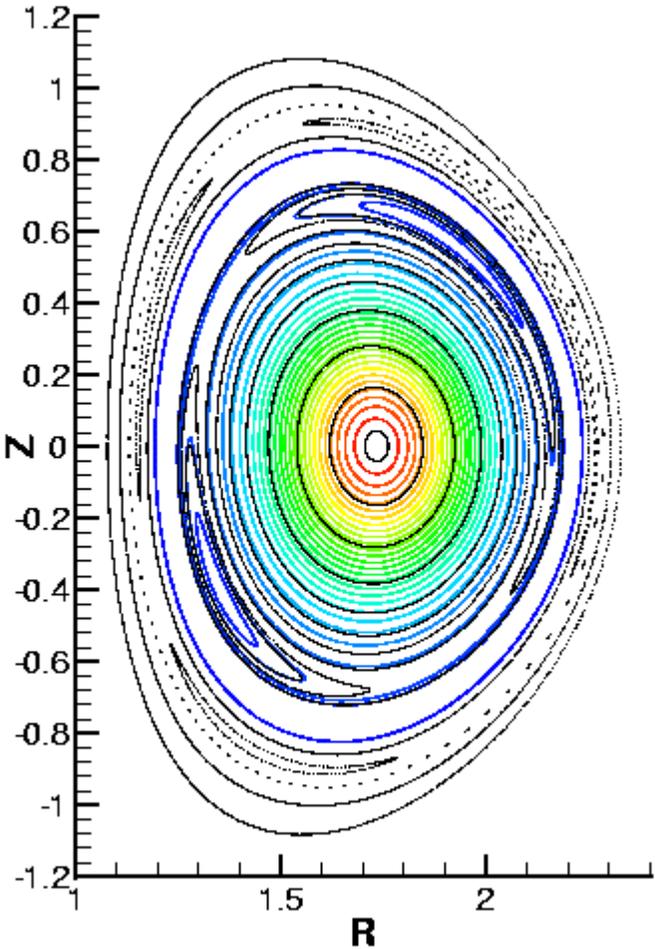
High-order basis functions enable resolution of extreme anisotropy in nonlinear simulations of electromagnetic fusion physics.

Finite elements also provide geometric flexibility.

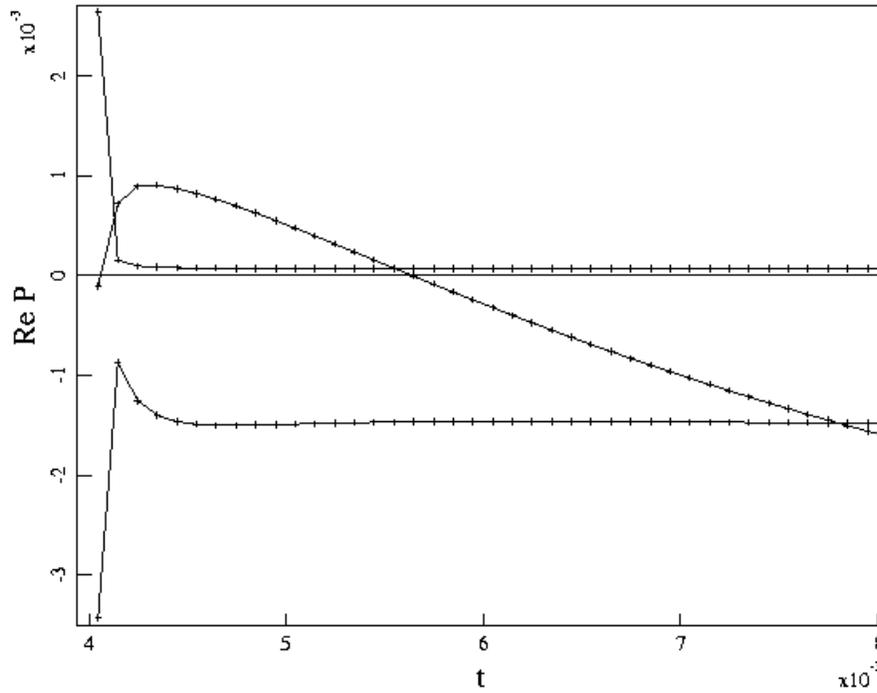
Anisotropic Diffusion Demonstration

- Start from $S=10^4$, $P_m=0.1$ saturation of Test Problem 1b (DIII-D-like equilibrium) that had $\kappa_{iso}=0.423 \text{ m}^2/\text{s}=0.1 \nu$.
- Freeze magnetic evolution and just run anisotropic thermal diffusion over the perpendicular time-scale (100 x resistive time-scale).
- $\kappa_{\perp}=0.423$, $\kappa_{\parallel}=4.23 \times 10^8$, $\Delta t=1 \times 10^{-4} \text{ s}$.
- Island appears in pressure contours immediately.

Pressure contours (color) after 1.5 ms of anisotropic diffusion overlaid with Poincare surface of section.

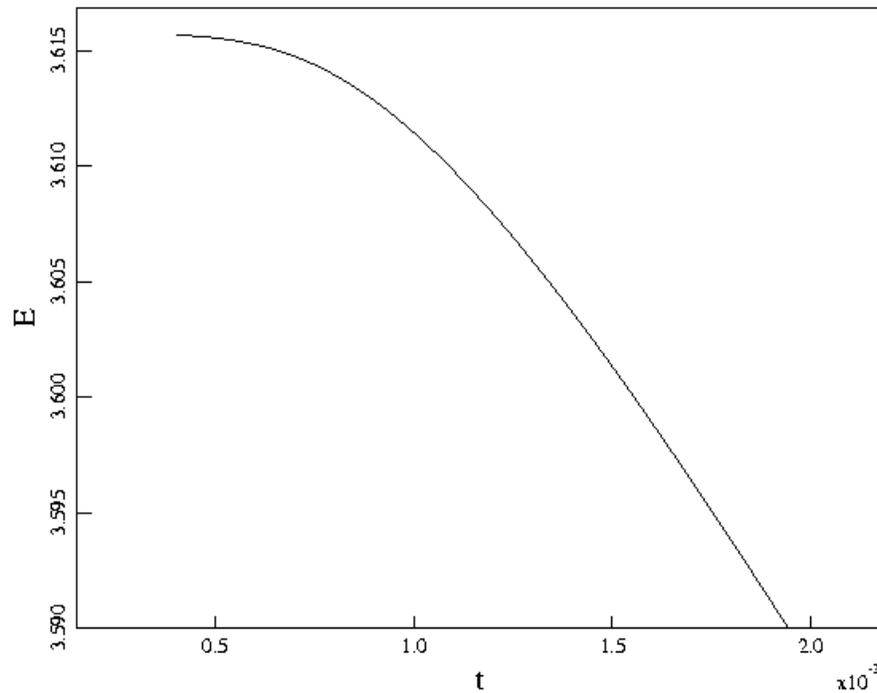


Re P vs. t



a)

Internal Energy vs. t

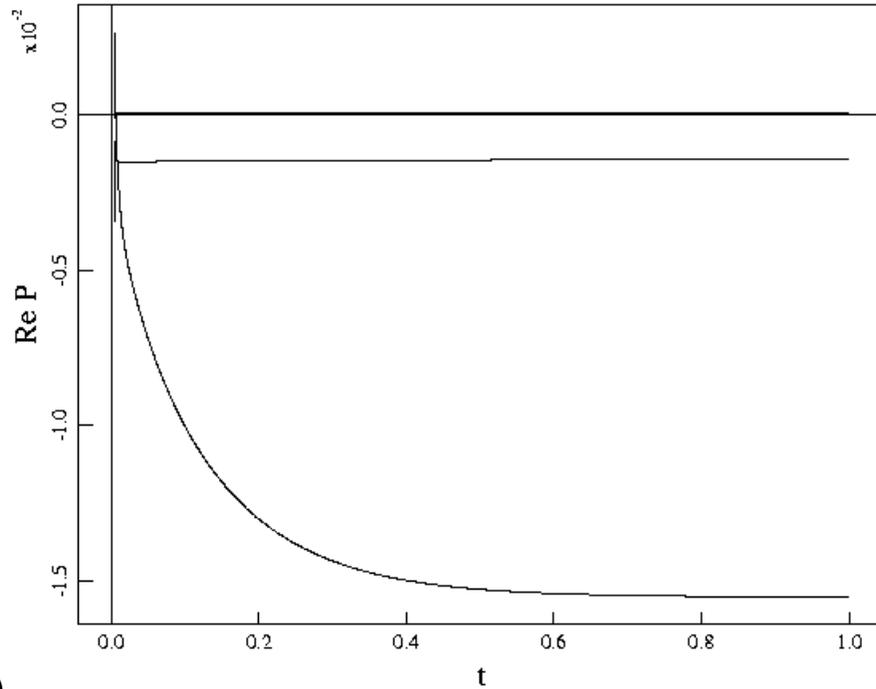


b)

a) "Probe" located at (R=1.312, Z=-0.248) and b) internal energy vs. time. Time-scale in a) is ms, and in b) it's 10 ms.

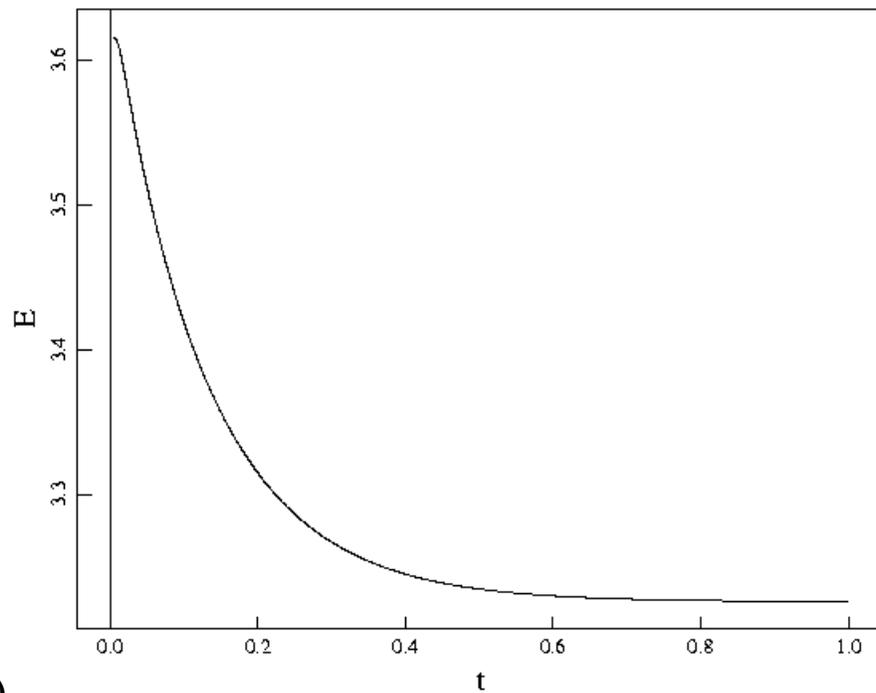
- Steady-state transport is regained after a perpendicular diffusion time.

Re P vs. t



a)

Internal Energy vs. t



b)

a) "Probe" and b) internal energy. Time-scale is seconds.

- The $n=0$ pressure profile inside the island drops over the long time-scale, reflecting the loss of insulation over the magnetic island.

